



Colgate-Palmolive

# Lattice QCD Input to Axion Cosmology

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Evan Berkowitz  
Lawrence Livermore National Laboratory

Workshop on Microwave Cavity Design for Axion Detection  
LLNL

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# Axions

Peccei & Quinn: PRL **38** (1977) 1440, PR **D16** (1977) 1791

- Couple to QCD topological charge

$$\mathcal{L}_{\text{axions}} = \frac{1}{2} (\partial_\mu a)^2 + \left( \frac{a}{f_a} + \theta \right) \frac{1}{32\pi^2} \epsilon^{\mu\nu\rho\sigma} F_{\mu\nu} F_{\rho\sigma}$$

- Mass is temperature dependent.

$$m_a^2 f_a^2 = \left. \frac{\partial^2 F}{\partial \theta^2} \right|_{\theta=0} \equiv \chi(T)$$

Axion mass    QCD Topological Susceptibility

- After production axions get diluted

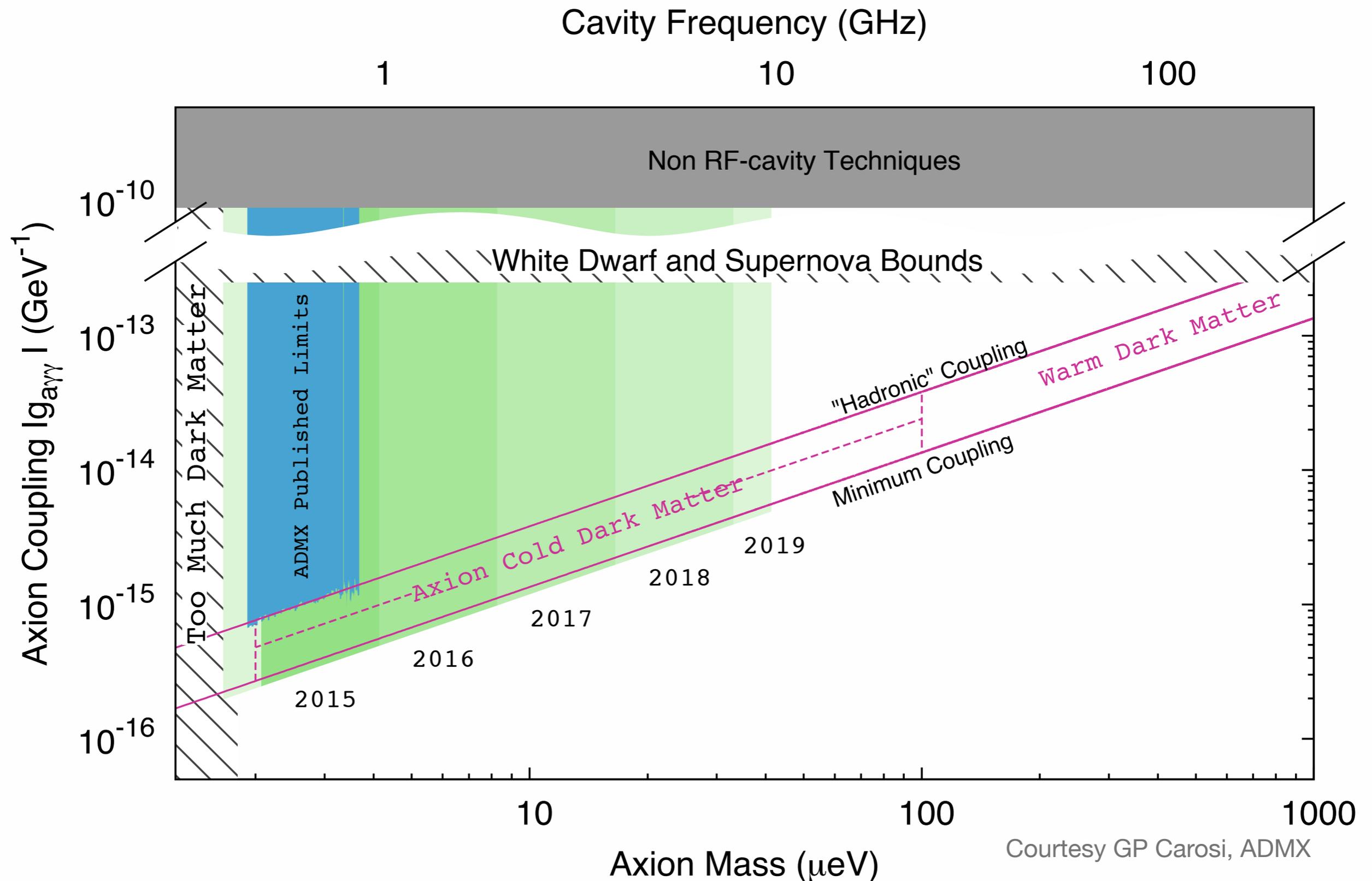
$$\frac{\rho(t) R^3}{m_a(t)} = \# \text{ axions in a fixed comoving volume}$$

- Requiring  $\rho(\text{today}) \leq \rho_{\text{DM}}$  gives an upper bound on  $f_a$  and a lower bound on  $m_a$  if PQ breaking is after inflation

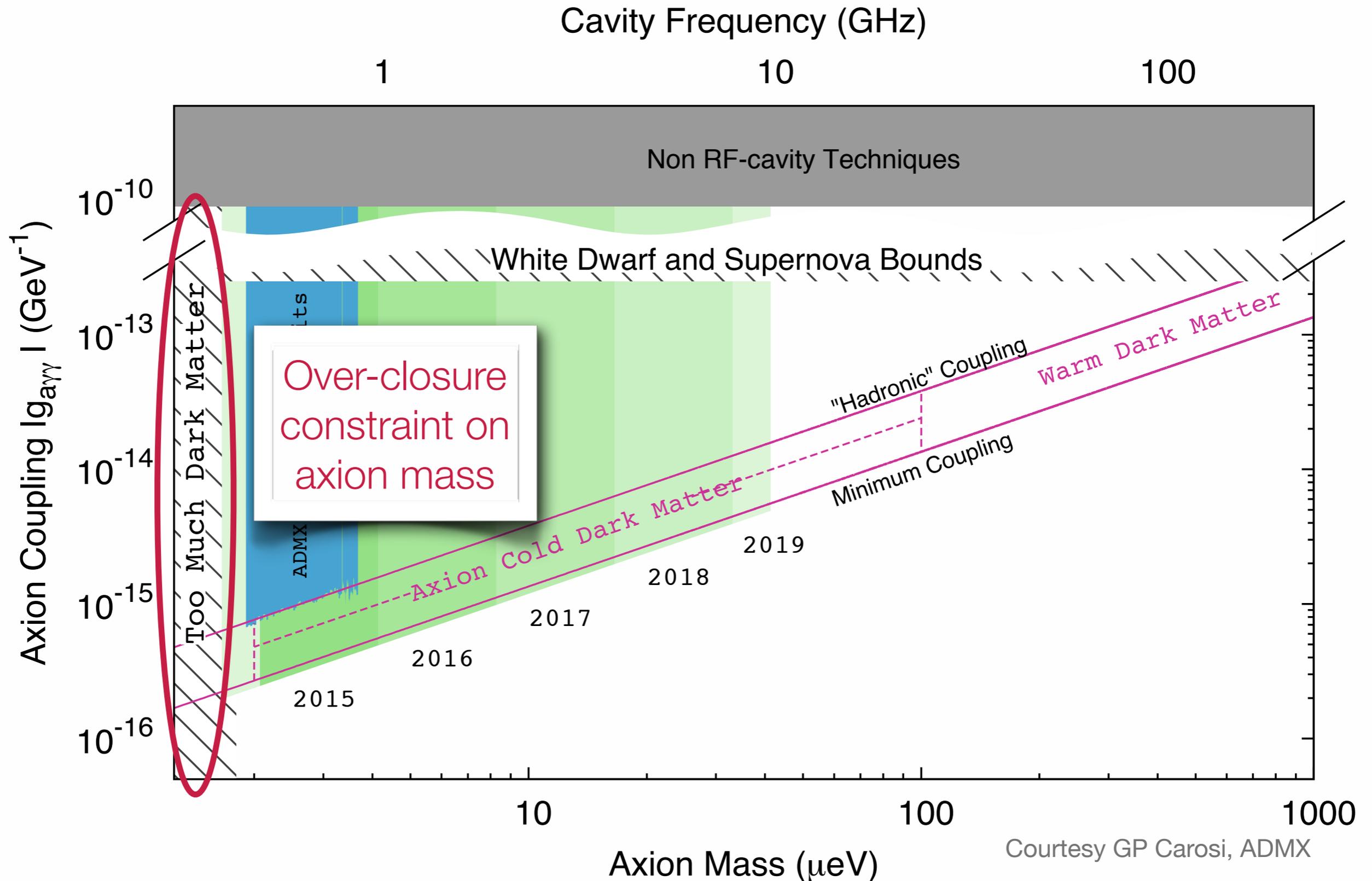
$$\rho(T_\gamma) = \frac{\pi^2}{6} \sqrt{\chi(T_\gamma) \chi(T_1)} \left( \frac{R(T_1)}{R(T_\gamma)} \right)^3$$

$\chi(T)$  ← our lattice QCD calculation →  $T_\gamma$  today  
 $\chi(T)$  ← our lattice QCD calculation →  $T_1$  production ceased  
 $\chi(T)$  ← our lattice QCD calculation →  $R$  GR scale factor

# Axion Constraints



# Axion Constraints



# CAVEAT

We currently study pure Yang-Mills (gluons only)  
and not yet full QCD!

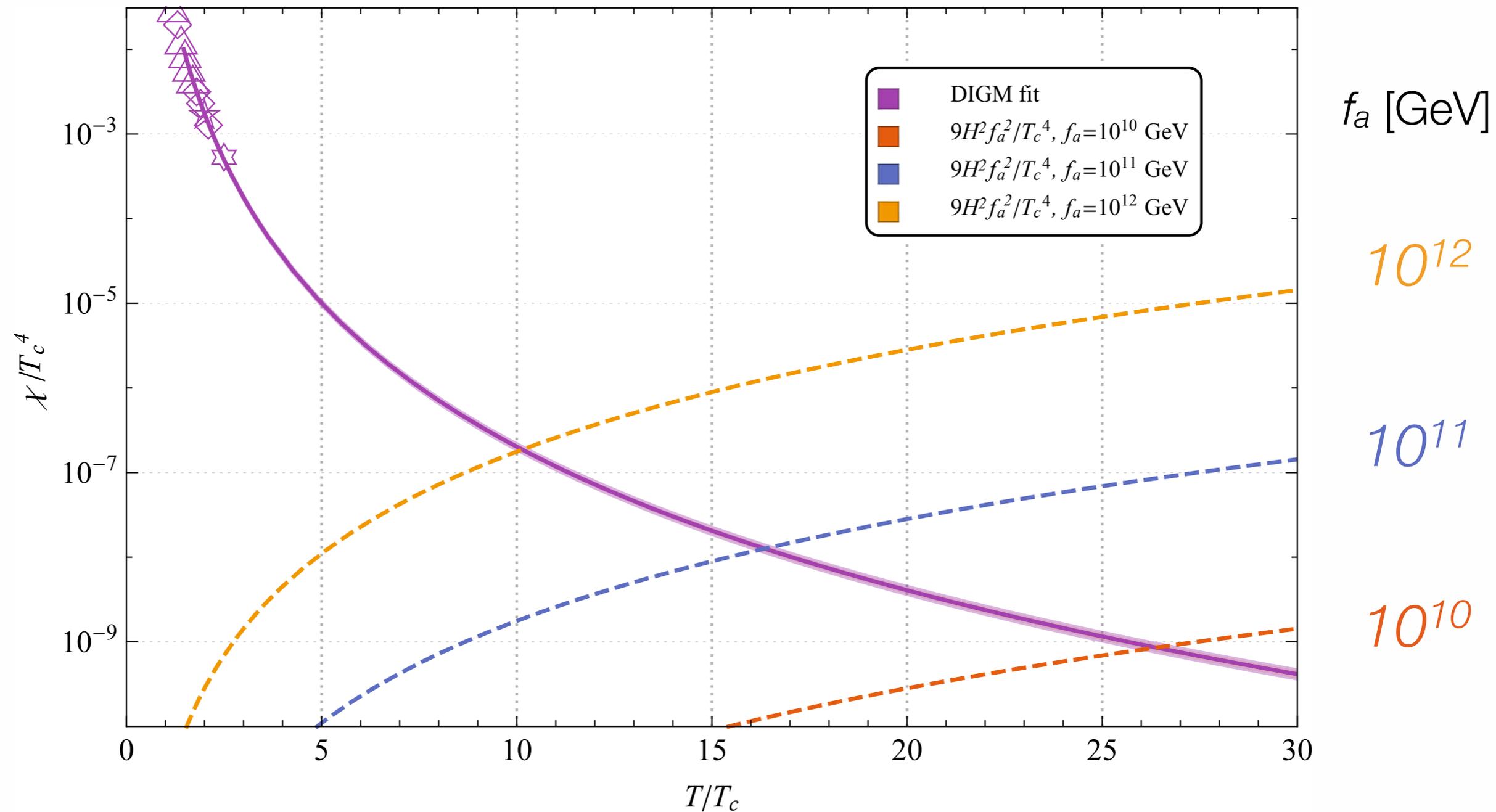
- Dramatically more efficient algorithms enable huge statistics and volumes, shorter autocorrelation times.
- Should agree at (asymptotically) high temperatures

But: for getting a lower bound on  $m_a$  it's not so bad!

- Pure glue is the  $\infty$  quark mass limit, and theory arguments imply  $\chi(T)$  should go to 0 when even one quark mass  $\rightarrow 0$ . So plausibly  $\chi(T)_{\text{QCD}} < \chi(T)_{\text{glue}}$
- Bound gets stronger when  $\chi(T)$  gets smaller.

# Axion production ceases when $9H^2 f_a^2 = m_a^2 f_a^2 = \chi$

Berkowitz, Buchoff & Rinaldi, PRD92.034507 arXiv:1505.07455



Intersection gives  $T_1$  as a function of  $f_a$



